

On the binary additive divisor problem in mean

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Abstract

We study a mean value of the classical additive divisor problem, that is

$$\sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l \sim L} d(n+l)d(n+l+f) - \text{main term} \right|^2,$$

with quantities $N \geq 1$, $1 \leq F \ll N^{1-\varepsilon}$ and $1 \leq L \leq N$. The main term we are interested in here is the one by Motohashi [27], but we also give an upper bound for the case where the main term is that of Atkinson [1]. Furthermore, we point out that the proof yields an analogous upper bound for a shifted convolution sum over Fourier coefficients of a fixed holomorphic cusp form in mean.

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1 Introduction

A widely studied occurrence of the classical divisor function $d(n) = \sum_{d|n} 1$ is in the additive divisor problem, in which one investigates the asymptotic behavior of the sum

$$D(x; m) = \sum_{n \leq x} d(n)d(n+m) \quad (1)$$

as x tends to infinity and m is a given positive integer. Ingham was the first to find an asymptotic formula for the above sum in 1927, and in 1931 Estermann brought the Kloosterman sums into discussion, improving the earlier result. The importance of uniformity in the shift m was noted ten years later by Atkinson [1], who further used his findings in studying the fourth power mean of the Riemann zeta-function on the critical line. In 1982 Deshouillers and Iwaniec applied Kuznetsov's trace formulas into the problem, obtaining a notable improvement to the upper bound of the error term. Finally, especially crucial for

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this paper, Motohashi [27] worked out an explicit spectral decomposition for a weighted convolution sum, separating a main term and obtaining an upper bound for the error term uniformly in the shift m . We mention shortly papers by Heath-Brown, Iwaniec [11] and Ivić-Motohashi [9], [10], who obtained mean value results for the error term, and [23], [25] and [32] for other related results. For a more thorough discussion of the history and the references of the above mentioned classical papers we refer to Motohashi's comprehensive paper [27].

In recent years the analogy between the divisor function and the Fourier coefficients of holomorphic and non-holomorphic cusp forms has appeared repeatedly in the literature. The behaviour of the analogous sums to (1),

$$A(x; m) = \sum_{n \leq x} a(n) \overline{a(n+m)} \quad (2)$$

involving the Fourier coefficients of a holomorphic cusp form and

$$T(x; m) = \sum_{n \leq x} t(n) t(n+m) \quad (3)$$

involving the Hecke eigenvalues corresponding to Fourier coefficients of a Maass form, has been studied intensively. In 1965 Selberg studied the meromorphy of a certain Dirichlet series related to the Fourier coefficients of holomorphic cusp forms, while Good in 1981 estimated the second moment of a modular L -function on the critical line using the spectral decomposition of shifted convolution sums coming from the Fourier coefficients of a holomorphic cusp form. Jutila [14], [15] derived explicit asymptotic formulae for all three cases (1), (2) and (3) in a unified way, using the respective generating Dirichlet series. For a discussion of the substantial history including e.g. the works by Blomer, Duke, Friedlander, Harcos, Iwaniec, Jutila, Michel, Motohashi and Sarnak as well as a general spectral decomposition for the weighted shifted convolution sums we refer to the paper [2] by Blomer and Harcos. Along with these references we mention also [3], [6], [7], [16] and [20]. Together the three sums above go under the name of the shifted convolution problem.

As a prologue for this paper, in Lemma 3 of his paper [16] Jutila studies the sum (2) in mean, and his argument can easily be extended to prove

$$\sum_{0 \leq f \leq F} \sum_{1 \leq n \leq N} \left| \sum_{1 \leq l \leq L} a(n+l) \overline{a(n+f+l)} \right|^2 \ll (N+F)^k N^k L \quad (4)$$

for all $N, F \geq 1$ and $1 \leq L \leq N$. Instantaneously the proof gives also an upper bound for a sum over Hecke eigenvalues of Maass forms in mean, as stated in Lemma 6 of [17] : For $N, F \geq 1$ and $1 \leq L \leq N$

$$\sum_{0 \leq f \leq F} \sum_{1 \leq n \leq N} \left| \sum_{1 \leq l \leq L} t(n+l) t(n+f+l) \right|^2 \ll (N+F)^{1+\varepsilon} N L. \quad (5)$$

The idea in the estimation of the triple sum lies in its sensitiveness of the size of the innermost sum over l . A similar result, sensitive for the size of the shift,

is needed in the doctoral thesis of the author ([31], Lemma 3.4). An analogous result over more general settings also appears on p. 81 in the paper [4] by Blomer, Harcos and Michel, leading to the same bounds (4), (5). In all the three papers [16], [17], [4] the motivation behind these studies has been gaining information about the upper bounds of the original shifted convolution sums (2) and (3), while the author needed her estimate in her thesis in estimating a certain spectral sum over inner products involving holomorphic cusp forms and Maass forms.

However, to our knowledge there does not exist a similar bound for the case of the divisor function, and the extension of the proofs of the cuspidal cases seem problematic due to the main term of the additive divisor problem. This paper is our attempt to bring new light for this case. Hence we state the following Theorem:

Theorem 1.1 (Main theorem). *Let $N \geq 1$, $1 \leq L \leq N$ and $1 \leq F \ll N^{1-\varepsilon}$. Then*

$$S = \sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l \sim L} d(n+l)d(n+l+f) - \frac{6}{\pi^2} \int_{(L+n)/f}^{(2L+n)/f} m(x; f) dx \right|^2$$

$$\ll N^{2+\varepsilon} + N^{1+\varepsilon} L F. \quad (6)$$

Here $m(x; f)$ is as in (1.12) in [27].

The notation $m \sim M$ stands for $M < m \leq 2M$. For the classical main term $M(N, f)$ by Atkinson ([1], p. 185) we easily obtain the following upper bound:

Corollary 1.2. *Let $N \geq 1$, $1 \leq L \leq N$ and $1 \leq F \ll N^{1-\varepsilon}$. Then*

$$S = \sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l \sim L} d(n+l)d(n+l+f) - M(2L+n, f) + M(L+n, f) \right|^2$$

$$\ll N^{2+\varepsilon} + N^{1+\varepsilon} L F + N^{-1+\varepsilon} L^2 F^3.$$

Remark 1.3. With small values of the quantities F and L , that is, if $L^2 F = o(N)$, the trivial bound $N^{1+\varepsilon} L^2 F$ yields a better result than Theorem 1.1 and Corollary 1.2. On the other hand, in their paper [10] Ivić and Motohashi gain the upper bound

$$\sum_{f=1}^F \left| \sum_{n=1}^N d(n)d(n+f) - \frac{6}{\pi^2} \int_0^{N/f} m(x; f) dx \right|^2 \ll N^{4/3+\varepsilon} F^{1/3}$$

uniformly for $F \leq N^{1/2-\varepsilon}$. Here, according to Meurman ([25], pp. 225-226) it seems possible to extend the range to $F \leq N$ by replacing the upper bound by $N^{4/3+\varepsilon} F^{1/3} + NF$. This result yields, up to factor N^ε , the same upper bound as in Theorem 1.1 in case $F \gg N^{1/2}$ and $L = N$. Other than this, to our knowledge, our result (6) and the trivial bound together yield the genuinely best upper bound with the restriction $F \ll N^{1-\varepsilon}$.

The main novelty in this paper is the methodological uniformity allowing an analogous treatment for the cuspidal cases: Following the analogous proof for the holomorphic cusp forms we immediately gain

Theorem 1.4. *Let $N \geq 1$, $1 \leq L \leq N$ and $1 \leq F \ll N^{1-\varepsilon}$. Then*

$$\sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l \sim L} a(n+l) \overline{a(n+l+f)} \right|^2 \ll N^{2k+\varepsilon} + N^{2k-1+\varepsilon} LF.$$

To compare our result to the earlier mentioned, we notice that up to a factor N^ε and with the restriction $F \ll N^{1-\varepsilon}$, (4) yields our result only if $L \asymp 1$, and otherwise our result is genuinely better.

In the case of non-holomorphic cusp forms we face the problem of the lack of a proper analogy for the spectral decomposition of the shifted convolution sum in question. There exists an analogous result for our crucial Lemma 2.3, namely Lemma 5 in [22], originally due to Jutila [13], [19]. However, although the analogous part of this decomposition results the anticipated analogous bound, there appears also an arithmetic correction term, which leads to additional problems. If we were dealing with a genuinely oscillating weight function in our proof, then the contribution of the correction term would easily be proven to stay under the expected bound. However, as this is not the case we shall lighten this paper and just assign the conjecture below, leaving the proof for future study.

Conjecture 1.5. *Let $N \geq 1$, $1 \leq L \leq N$ and $1 \leq F \ll N^{1-\varepsilon}$. Then*

$$\sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l \sim L} t(n+l) t(n+l+f) \right|^2 \ll N^{2+\varepsilon} + N^{1+\varepsilon} LF.$$

Remark 1.6. Note also Theorem 1 in [2], which gives a more general spectral decomposition of a shifted convolution sum including Hecke eigenvalues of the underlying cusp forms of two arbitrary cuspidal automorphic representations of $PGL_2(\mathbb{Z}) \backslash PGL_2(\mathbb{R})$.

2 The needed notation and auxiliary lemmas

2.1 The cusp forms

Besides their analogous behavior to the divisor function, the holomorphic and non-holomorphic cusp forms play a role in the spectral decomposition of the shifted convolution sum over the divisor function (Lemma 2.3). Hence we start by introducing briefly some results concerning them. For the proofs and for a general reference the reader is referred to Motohashi's monograph [29].

We may confine ourselves to the cusp forms for the full modular group $\Gamma = SL_2(\mathbb{Z})$ operating through Möbius transformations on the upper half plane \mathbb{H} .

A holomorphic cusp form $F(z) : \mathbb{H} \rightarrow \mathbb{C}$ of weight $k \in \mathbb{Z}$ with respect to Γ can be represented by its Fourier series

$$F(z) = \sum_{n=1}^{\infty} a(n)e(nz), \quad e(\alpha) = \exp(2\pi i\alpha).$$

We may assume that k is even and $k \geq 12$, otherwise $F(z)$ is trivial. We let

$$\{\psi_{j,k} \mid 1 \leq j \leq \vartheta(k)\}$$

be an orthonormal basis of the unitary space of holomorphic cusp forms of weight k , and write

$$\psi_{j,k}(z) = \sum_{n=1}^{\infty} \rho_{j,k}(n) n^{\frac{k-1}{2}} e(nz).$$

We may suppose that the basis vectors are eigenfunctions of the Hecke operators $T_k(n)$ for all positive integers n . Thus, in particular, $T_k(n)\psi_{j,k} = t_{j,k}(n)\psi_{j,k}$ for certain real numbers $t_{j,k}(n)$, which we call Hecke eigenvalues. Comparing Fourier coefficients on both sides, one may verify that $\rho_{j,k}(n) = \rho_{j,k}(1)t_{j,k}(n)$ for all $n \geq 1$, $1 \leq j \leq \vartheta(k)$. We put $\alpha_{j,k} = 16\Gamma(2k)(4\pi)^{-2k-1} \times |\rho_{j,k}(1)|^2$, whence

$$\sum_{j=1}^{\vartheta(k)} \alpha_{j,k} \ll k. \quad (7)$$

Furthermore, we let

$$H_{j,k}(s) = \sum_{n=1}^{\infty} t_{j,k}(n)n^{-s}$$

stand for the Hecke L-function attached to $\psi_{j,k}$. The series converges absolutely for $\text{Re } s > 1$ because of the bound

$$|t_{j,k}(n)| \leq d(n) \ll n^{\varepsilon} \quad (8)$$

by Deligne [5]. Furthermore $H_{j,k}(s)$ can be continued to an entire function, and it satisfies a functional equation, which implies that for bounded s

$$H_{j,k}(s) \ll k^c \quad (9)$$

uniformly in j . Here c is a suitable constant, which depends only on $\text{Re } s$.

A non-holomorphic cusp form $u(z) = u(x + iy) : \mathbb{H} \rightarrow \mathbb{C}$ is a non-constant real-analytic Γ -invariant function in the upper half-plane, square-integrable with respect to the hyperbolic measure $d\mu(z) = \frac{dx dy}{y^2}$ over a fundamental domain of Γ . Furthermore $u(z)$ is an eigenfunction of the non-euclidean Laplacian $\Delta = -y^2(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2})$, and the corresponding eigenvalue can be written as $1/4 + \kappa^2$ with $\kappa > 0$. The Fourier series expansion for $u(z)$ is then of the form

$$u(z) = y^{1/2} \sum_{n \neq 0} \rho(n) K_{i\kappa}(2\pi|n|y) e(nx)$$

with K_ν a Bessel function of imaginary argument. We may suppose that our cusp forms are eigenfunctions of the Hecke operators $T(n)$ for all positive integers n and that $u(x+iy)$ is even or odd as a function of x . Thus $T(n)u = t(n)u$ for certain real numbers $t(n)$, which are again called Hecke eigenvalues, and $u(-\bar{z}) = \pm u(z)$. Comparing Fourier coefficients on both sides, one may verify that $\rho(n) = \rho(1)t(n)$ and $\rho(-n) = \pm \rho(n)$ for all $n \geq 1$. The *Maass (wave) forms* u_j constitute an orthonormal set of non-holomorphic cusp forms arranged so that the corresponding parameters κ_j determined by the eigenvalues $1/4 + \kappa_j^2$ lie in an increasing order. We write $\rho_j(n)$ and $t_j(n)$ for the corresponding Fourier coefficients and Hecke eigenvalues. Now let

$$\alpha_j = |\rho_j(1)|^2 / \cosh(\pi\kappa_j),$$

and let

$$H_j(s) = \sum_{n=1}^{\infty} t_j(n)n^{-s}$$

stand for the Hecke L -function attached to u_j . As the Ramanujan-Petersson conjecture $t_j(n) \ll n^\varepsilon$ holds in a mean value sense, the series converges absolutely for $\operatorname{Re} s > 1$. Again, $H_j(s)$ can be continued to an entire function. We have the bound

$$\sum_{\kappa_j \leq K} \alpha_j H_j^4\left(\frac{1}{2}\right) \ll K^2 \log^{15} K. \quad (10)$$

(See [29], Theorem 3.4.)

Lastly, in this paper, the following notation will be adopted: Vinogradov's relation $f(z) \ll g(z)$ is another notation for $f(z) = \mathcal{O}(g(z))$. We let ε stand generally for a small positive number, not necessarily the same at each occurrence.

2.2 The needed lemmas

In this section we shall gather some auxiliary results which will be used during the course of the proofs of the theorems.

For estimating the Gamma function we shall frequently use the following Stirling's formula without further reference: In any fixed strip $b \leq \sigma \leq c$ we have

$$|\Gamma(\sigma + it)| = \sqrt{2\pi} |t|^{\sigma-1/2} e^{-|t|\pi/2} (1 + \mathcal{O}(|t|^{-1})),$$

for $|t| \rightarrow \infty$. For a proof, see Olver [30] p. 294. Further we have

$$\Gamma(x) = \sqrt{2\pi} x^{x-1/2} e^{-x} (1 + \mathcal{O}(x^{-1})), \quad (11)$$

for all positive real $x \rightarrow \infty$, see [24], Eq. (1.4.25).

The proof of the following classical estimate for the fourth moment of Riemann's zeta-function on the critical line can be found e.g. from [33], p. 147:

$$\int_0^T \left| \zeta\left(\frac{1}{2} + it\right) \right|^4 dt \ll T \log^4 T \quad (12)$$

for $T \geq 1$. Along with this estimate we have

$$\int_0^T \left| \zeta \left(\frac{1}{2} + it \right) \right|^{12} dt \ll T^2 \log^{17} T \quad (13)$$

for $T \geq 1$ by Heath-Brown [8]. Moreover it is known that

$$\frac{1}{\zeta(1+it)} \ll \log |t|,$$

as $|t| \geq 1$. For a proof, see e.g. [33], p. 132.

We have an important tool arising from spectral theory:

Lemma 2.1 (The spectral large sieve). *For $K \geq 1$, $1 \leq \Delta \leq K$, $M \geq 1$ and any complex numbers a_m we have*

$$\sum_{K \leq \kappa_j \leq K+\Delta} \alpha_j \left| \sum_{m \leq M} a_m t_j(m) \right|^2 \ll (K\Delta + M)(KM)^\varepsilon \sum_{m \leq M} |a_m|^2.$$

For a proof, see Theorem 1.1 in [18] or Theorem 3.3 in [29]. The continuous analogy is the following:

Lemma 2.2. *For K real, $\Delta \geq 0$, $M \geq 1$ and any complex numbers a_m we have*

$$\int_K^{K+\Delta} \left| \sum_{m \leq M} a_m \sigma_{2ir}(m) m^{-ir} \right|^2 dr \ll (\Delta^2 + M) M^\varepsilon \sum_{m \leq M} |a_m|^2$$

uniformly in K , as $\sigma_{2ir}(f) = \sum_{d|f} d^{2ir}$.

Proof can be found in [31], Lemma 1.12.

A crucial role is played by a spectral decomposition of the shifted convolution sum over the divisor function:

Lemma 2.3. *Let f be a positive integer and W a smooth function of compact support on $(0, \infty)$. Then*

$$\begin{aligned} \sum_{l=1}^{\infty} d(l) d(l+f) W\left(\frac{l}{f}\right) &= \frac{6}{\pi^2} \int_0^{\infty} m(x; f) W(x) dx + \frac{f^{1/2}}{\pi} \int_{-\infty}^{\infty} f^{-ir} \sigma_{2ir}(f) \\ &\times \frac{|\zeta(\frac{1}{2} + ir)|^4}{|\zeta(1 + 2ir)|^2} \Theta(r; W) dr + f^{1/2} \sum_{j=1}^{\infty} \alpha_j t_j(f) H_j^2\left(\frac{1}{2}\right) \Theta(\kappa_j; W) \\ &+ \frac{1}{4} f^{1/2} \sum_{k=6}^{\infty} \sum_{j=1}^{\vartheta(k)} (-1)^k \alpha_{j,k} t_{j,k}(f) H_{j,k}^2\left(\frac{1}{2}\right) \Xi\left(k - \frac{1}{2}; W\right) \end{aligned} \quad (14)$$

Here $m(x; f)$ is defined by (1.12) in [27], $\sigma_{2ir}(f)$ is as in the previous lemma and

$$\begin{aligned}\Theta(r; W) &= \frac{1}{2} \operatorname{Re} \left\{ \left(1 + \frac{i}{\sinh(\pi r)} \right) \Xi(ir; W) \right\} \\ &= \frac{1}{4} \left(1 + \frac{i}{\sinh(\pi r)} \right) \Xi(ir; W) + (r \mapsto -r),\end{aligned}\tag{15}$$

with

$$\begin{aligned}\Xi(\xi; W) &= \frac{\Gamma(\xi + \frac{1}{2})^2}{\Gamma(2\xi + 1)} \int_0^\infty x^{-1/2-\xi} F\left(\xi + \frac{1}{2}, \xi + \frac{1}{2}; 2\xi + 1; -\frac{1}{x}\right) \\ &\quad \times W(x) dx.\end{aligned}\tag{16}$$

Here $F(., .; .; .)$ is the hypergeometric function and $(r \mapsto -r)$ stands for an expression similar to the preceding one, but with r replaced by $-r$.

For the proof, see [27], Theorem 3.

Finally we introduce the basic inequality in the proof of the classical large sieve:

Lemma 2.4 (Sobolev). *Let $a \leq u \leq a + \Delta$ and let the function f be continuously differentiable on this interval. Then*

$$\begin{aligned}|f(u)|^2 &\leq \Delta^{-1} \int_a^{a+\Delta} |f(x)|^2 dx + 2 \left(\int_a^{a+\Delta} |f(x)|^2 dx \right)^{1/2} \\ &\quad \times \left(\int_a^{a+\Delta} |f'(x)|^2 dx \right)^{1/2} \ll \Delta^{-1} \int_a^{a+\Delta} (|f(x)|^2 + \Delta^2 |f'(x)|^2) dx\end{aligned}$$

uniformly.

For a proof, see Montgomery [26], Lemma 1.1 applied to f^2 .

3 Proofs of the results

3.1 Theorem 1.1

We start by adding a set of $\asymp \log(1/U)$ real-valued smooth weight functions $g_\delta(\frac{l}{L})$ to the l -sum so that their sum produces an approximation of the characteristic function of the interval $[1, 2]$ with an error of size $\ll U$ and their supports widen step by step by factors 2 when we move away from the end points 1 and 2. Thus if the supports of the first and last weight functions are of length $U = L^{-1+\varepsilon}$, then the next ones are of length $\asymp 2U$ and so on, and the slopes of the weight functions cancel out each other. Hence we let $g_\delta(x)$ stand for a compactly supported function on some interval of length $\asymp \delta$ contained in

[1, 2]. Moreover, $g(x) = 1$ on an interval of length $\asymp \delta$ and $g^{(\nu)}(x) \ll_{\nu} \delta^{-\nu}$ for each $\nu \geq 0$ and $x \in \mathbb{R}$. Here $L^{-1+\varepsilon} \leq \delta \ll 1$. Now

$$\begin{aligned} S &\ll L^{\varepsilon} \sum_{f \sim F} \sum_{n \sim N} \left| \sum_{l=1}^{\infty} d(n+l) d(n+l+f) W\left(\frac{n+l}{f}\right) \right. \\ &\quad \left. - \frac{6}{\pi^2} \int_0^{\infty} m(x; f) W(x) dx \right|^2 + N^{1+\varepsilon} F U^2 L^2 \\ &\ll L^{\varepsilon} \sum_{f \sim F} \sum_{n \sim N} |e_1(n, f; L) + e_2(n, f; L) + e_3(n, f; L)|^2 + N^{1+\varepsilon} F, \end{aligned}$$

with $W(x) = g_{\delta}\left(\frac{fx-n}{L}\right)$. Here $e_{\nu}(n, f; L)$ ($\nu = 1, 2, 3$) stands for the $(\nu + 1)$ th term on the right hand side of (14).

The contribution of the term $e_2(n, f; L)$ turns out to be the most difficult to bound. Let us denote

$$S_2 = \sum_{f \sim F} \sum_{n \sim N} \left| f^{1/2} \sum_{j=1}^{\infty} \alpha_j t_j(f) H_j^2\left(\frac{1}{2}\right) \Theta(\kappa_j; W) \right|^2$$

with $\Theta(r; W)$ as in (15). As the support of W tends to infinity, we have in the integral in (16)

$$\begin{aligned} F \left(\frac{1}{2} + ir, \frac{1}{2} + ir; 1 + 2ir; -\frac{1}{x} \right) &= \left(\frac{1 + \sqrt{1 + x^{-1}}}{2} \right)^{-1-2ir} \\ &\times F \left(\frac{1}{2} + ir, \frac{1}{2}; 1 + ir; \left(\frac{1 - \sqrt{1 + x^{-1}}}{1 + \sqrt{1 + x^{-1}}} \right)^2 \right) \end{aligned} \quad (17)$$

by the quadratic transformation formula for the hypergeometric function from Lebedev [24], Eq. (9.6.12). We write the RHS of (17) by the definition of the hypergeometric series. Now clearly the k th term in the series is non-oscillating and of order x^{-2k} . Therefore it suffices to study only the leading term 1, others behaving similarly until some sufficient constant, and the contribution of the rest being negligibly small. (See also Convention 2 in [21].)

By integrating repeatedly by parts we notice that κ_j can be truncated to be $\ll N^{1+\varepsilon} \delta^{-1} L^{-1}$ with negligibly small error. We further decompose the range $[1, N^{1+\varepsilon} \delta^{-1} L^{-1}]$ into dyadic intervals $\kappa_j \sim K$ with $1 \leq K \ll N^{1+\varepsilon} \delta^{-1} L^{-1}$, the number of these intervals being $\ll N^{\varepsilon}$. Next we prepare ourselves to use the duality principle (see for example pp. 169-170 in [12]). Let $b_{f,n}$ stand for any finite sequence of complex numbers such that $\sum_{f \sim F} \sum_{n \sim N} |b_{f,n}|^2 = 1$. Then

$$S_2 \ll N^{\varepsilon} L^2 F^{-1} \sum_{\kappa_j \sim K} \alpha_j \left| \sum_{f \sim F} \sum_{n \sim N} b_{f,n} t_j(f) I(\kappa_j; W) \right|^2 \sum_{\kappa_j \sim K} \alpha_j H_j^4\left(\frac{1}{2}\right) \kappa_j^{-1}$$

with

$$I(\kappa_j; W) = \int_1^2 \left(\frac{zL+n}{f} \right)^{-1/2-i\kappa_j} \left(\frac{1 + \sqrt{1 + \frac{f}{zL+n}}}{2} \right)^{-1-2i\kappa_j} g_\delta(z) dz.$$

Note that κ_j replaced by $-\kappa_j$ can be taken care of with complex conjugation. By (10) this is further

$$\ll N^\varepsilon K L^2 F^{-1} \sum_{\kappa_j \sim K} \alpha_j \left| \sum_{f \sim F} \sum_{n \sim N} b_{f,n} t_j(f) I(\kappa_j; W) \right|^2. \quad (18)$$

We want to next apply the spectral large sieve, so we shall use Sobolev's lemma 2.4 as done in Jutila's paper [14] on p. 454 to get rid of the parameter κ_j in the integral $I(\kappa_j; W)$. Therefore the range $(K, 2K]$ for κ_j is split up into segments of length Δ in such a way that the integral $I(y; W)$ remains essentially stationary as y runs over a segment. That is

$$\Delta \log \left(\frac{(zL+n) \left(1 + \sqrt{1 + \frac{f}{zL+n}} \right)^2}{4f} \right) \ll \log N.$$

In this way, the second term in the upper bound in Lemma 2.4 will be comparable to the first. An appropriate choice would be $\Delta = 1$.

Now we divide the κ_j -sum in (18) into subsums of length Δ , and apply Lemma 2.4 to each subsum arriving at the bound

$$\begin{aligned} N^\varepsilon K L^2 F^{-1} \sum_{l=0}^{K-1} \int_{K+l}^{K+l+1} \sum_{K+l < \kappa_j \leq K+l+1} \alpha_j \left(\left| \sum_{f \sim F} \sum_{n \sim N} b_{f,n} t_j(f) I(y; W) \right|^2 \right. \\ \left. + \left| \sum_{f \sim F} \sum_{n \sim N} b_{f,n} t_j(f) \frac{\partial}{\partial y} I(y; W) \right|^2 \right) dy. \end{aligned}$$

We next apply the spectral large sieve to the subsum over κ_j , and finally add the results corresponding to all subsums. This leads us to the bound

$$\begin{aligned} N^\varepsilon K L^2 F^{-1} (K+F) \sum_{f \sim F} \int_K^{2K} \left(\left| \sum_{n \sim N} b_{f,n} I(y; W) \right|^2 \right. \\ \left. + \left| \sum_{n \sim N} b_{f,n} \frac{\partial}{\partial y} I(y; W) \right|^2 \right) dy. \end{aligned}$$

Now we use Cauchy's inequality for the integrals $I(y; W)$ and $\frac{\partial}{\partial y} I(y; W)$, and attach a suitable real-valued smooth weight function $u(y)$ to the y -integral. We set u to be compactly supported on the interval $[K/2, 3K]$, $u(y) = 1$ on $[K, 2K]$

and $u^{(\nu)}(y) \ll_{\nu} K^{-\nu}$ for each $\nu \geq 0$ and $y \in \mathbb{R}$. By opening the squares we end up with a double sum $\sum_{n_1 \sim N} \sum_{n_2 \sim N}$. By repeated partial integration over the y -integral we now see that we may truncate $n_1 - n_2 \ll N^{1+\varepsilon} K^{-1}$. Hence we finally have an upper bound

$$N^{-1+\varepsilon} K^2 L^2 (K + F) \delta^2 \sum_{f \sim F} \sum_{n \sim N} |b_{f,n}|^2 \sum_{|n-n'| \ll N^{1+\varepsilon} K^{-1}} 1,$$

and we conclude with the desired upper bound.

The contribution of $e_1(n, f; L)$ to S can be estimated by exactly the same argumentation as above. We need the estimation

$$\int_0^T \left| \zeta \left(\frac{1}{2} + it \right) \right|^8 dt \ll T^{3/2+\varepsilon} \quad (19)$$

following readily of (12) and (13) by the Cauchy's inequality. Hence

$$S_1 \ll N^{1+\varepsilon} L F.$$

We could use also a more straightforward argumentation, but these steps make it easier to follow the proof of Theorem 1.4.

Using (7), (8), (9) and (11) the trivial estimates suffice for the estimation of $e_3(n, f; L)$.

3.2 Corollary 1.2

Using straightforward calculation it is easily seen that the difference of our main term and that of Atkinson's is

$$\frac{6}{\pi^2} \int_{(L+n)/f}^{(2L+n)/f} m(x; f) dx - (M(2L+n, f) - M(L+n, f)) \ll F L N^{-1+\varepsilon},$$

whence the result follows.

3.3 Theorem 1.4

Analogous to the proof of Theorem 1.1. For the size of the Fourier coefficients see (8). Instead of Lemma 2.3 we use Lemma 4 in [22], originally due to Motohashi [28]. The counterpart for (10) is Lemma 4 in [14], and instead of (19) we need Lemma 3 in [14].

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